



**Abstract**

Polarimetric coastal radar data are used to compare the rainfall characteristics of Hurricanes Harvey (2017) and Florence (2018). Intense rainfall was an infrequent yet important contributor to the total rainfall in Harvey, but its relative contribution varied spatially. The total rainfall over land maximized near the coast over Beaumont, TX due to intense convection resulting from prolonged onshore flow downshear from the circulation center. Overall, polarimetric radar observations in Harvey show a dominance of high concentrations of small-to-medium drops, consistent with prior tropical cyclone studies. The microphysical characteristics were spatially and temporally inhomogeneous however, with larger drops more frequent on 27 August and higher number concentrations more frequent on 28 and 30 August. The polarimetric variables and raindrop characteristics observed during Florence share broad similarities to Harvey, but had reduced variability, fewer observations of stronger reflectivity and differential reflectivity, and a lower frequency of high number concentrations and medium-sized drops. The radar data indicate Florence had reduced coverage of stronger convection compared to Harvey. We hypothesize that differences in storm motion, intensity decay rates, and vertical wind shear produce the distinct precipitation structures and microphysical differences seen in Harvey and Florence.

**1 Introduction**

Heavy rainfall is a known hazard of tropical cyclones (TCs), responsible for 25% of hurricane fatalities in the United States (*Rappaport* [2014]). During the 2017 and 2018 Atlantic hurricane seasons, rainfall from Hurricanes Harvey and Florence set new state rainfall records in Texas and the Carolinas and caused destructive flooding. Over a 6-d period in August 2017, Hurricane Harvey stalled over coastal Texas and the peak rainfall observation of 1538 mm near Nederland, TX broke the continental and overall United States TC rainfall records previously held by Tropical Storm Amelia (1978) and Hurricane Hiki (1950), respectively. Equally noteworthy, rainfall exceeded 500 mm over a large area extending from southeast of Austin, TX to the Texas-Louisiana border (*Blake and Zelinsky* [2018]). Just one year later, Hurricane Florence crept along the east coast of the United States, where accumulated rainfall totals of 912 mm over three days near Elizabethtown, NC and 600 mm near Loris, SC broke the TC precipitation records for each state. In each case, the extreme rainfall caused considerable impacts. Harvey caused over 65 direct deaths and widespread structural damage throughout southeast Texas (*Blake and*

45 *Zelinsky* [2018]). In the Carolinas, Florence caused 22 direct fatalities, widespread flood-  
46 ing, and damage (*Stewart and Berg* [2019]).

47 Forecasting extreme rainfall is challenging since the accumulation results from in-  
48 tense rain rates, long duration events, or a combination of both factors. The microphysical  
49 processes that determine the surface rainfall intensity and drop size distribution (DSD)  
50 in a TC are a complex mixture of raindrop creation, growth, and melting ice (*Black and*  
51 *Hallett* [1986]; *Marks and Houze* [1987]; *Houze et al.* [1992]; *Black and Hallett* [1999]).  
52 While the rotational storm dynamics and their interaction with the boundary layer are  
53 largely responsible for producing the vertical motion leading to precipitation, the envi-  
54 ronmental shear and the convective lifecycle also contribute to the strength and location  
55 of convective and stratiform precipitation in a TC (*Corbosiero and Molinari* [2002]; *Chen*  
56 *et al.* [2006]; *Hence and Houze* [2011, 2012]; *Didlake and Kumjian* [2017]). Overall, the  
57 complicated interplay between microphysical processes and the kinematic and thermo-  
58 dynamic environments that determine their importance, all within a translating, rotating  
59 storm, make it a challenge to accurately predict the local rainfall accumulation.

60 In situ and surface microphysical observations from rain gauges and disdrometers  
61 provide important information about hydrometeor characteristics, but only represent a  
62 small region of a TC. Nonetheless, disdrometer observations indicate that TC surface  
63 rainfall is typically dominated by numerous, midsize raindrops with median diameters  
64 of roughly 1-2 mm (*Jorgensen and Willis* [1982]; *Ulbrich and Lee* [2002]; *Tokay et al.*  
65 [2008]; *Chang et al.* [2009]; *Wang et al.* [2016]). DSD characteristics can vary substan-  
66 tially in time and space and depend on the specific precipitation growth mechanisms. In  
67 Typhoon Haima (2004), convective rainfall periods were composed of numerous midsize  
68 drops with median diameters around 2 mm, whereas stratiform rainfall periods had fewer  
69 drops with smaller median diameters between 1.25-1.75 mm (*Chang et al.* [2009]). How-  
70 ever, the wind speed limitations of ground-based disdrometers restrict most studies to re-  
71 gions away from the eyewall (*Ulbrich and Lee* [2002]; *Tokay et al.* [2008]; *Chang et al.*  
72 [2009]; *Wang et al.* [2016]).

73 Polarimetric radars can provide additional insight into the bulk hydrometeor shapes  
74 and concentrations over a much larger area (*Bringi and Chandrasekar* [2001]; *Ryzhkov*  
75 *et al.* [2005]; *Kumjian* [2013]; *Ryzhkov and Zrnica* [2019]). Although operational polarimet-  
76 ric radars in the United States are relatively recent, their data have improved our under-

77 standing of TC microphysical processes. *Brown et al.* [2016] analyzed the rainfall charac-  
78 teristics from Hurricanes Arthur and Ana (2014) and found similar but distinct probability  
79 distributions of the DSDs in the two TCs that were in some cases substantially different  
80 than the DSDs produced by numerical simulations. *Didlake and Kumjian* [2017] also an-  
81 alyzed Hurricane Arthur (2014) and found that convection located in the downshear right  
82 quadrant of the TC produced columnar and planar crystals, which were advected down-  
83 stream to the downshear left quadrant and fell as stratiform rain; despite different con-  
84 tributions from rain and ice processes, both quadrants exhibited heavy rain. The relative  
85 importance of rain and ice processes varies based on the precipitation feature and type.

86 Polarimetric analysis of the convective regions of a single rainband in Typhoon  
87 Matmo (2014) revealed a higher contribution of warm-rain processes to the surface rain,  
88 though ice processes were not insignificant (*Wang et al.* [2016]). A case study of an outer  
89 rainband in Typhoon Nida (2016) showed more influence from convective ice processes  
90 (*Wu et al.* [2018]). Case studies of Hurricanes Irene (2011) and Arthur (2014) revealed  
91 small ice was a weak contributor to the total ice water path in convective and stratiform  
92 precipitation, but was prevalent in low-reflectivity regions (*Kalina et al.* [2017]). In addi-  
93 tion to differences associated with specific features, processes can be influenced by envi-  
94 ronmental factors. Recently, *Didlake and Kumjian* [2018] and *Feng and Bell* [2019] found  
95 evidence of size-sorting of raindrops in TCs due to the asymmetric vertical motion caused  
96 by the storm motion and deep layer vertical wind shear vector, respectively. The variety of  
97 conclusions with regards to the importance of different microphysical and dynamical pro-  
98 cesses in TCs underscore the spatial and temporal variability of rainfall production mecha-  
99 nisms as well as the variability in storm structure and environmental conditions.

100 Hurricanes Harvey and Florence offer an opportunity to study the bulk surface rain-  
101 fall characteristics in two record-setting storms. *Wolff et al.* [2019] used polarimetric radar  
102 observations of Harvey to show that distinct DSD regimes impact rain rate retrievals and  
103 complicate the selection of parameters for attenuation-based algorithms. Through analyz-  
104 ing reanalysis and polarimetric radar data, *Brauer et al.* [2020] showed that strong hori-  
105 zontal moisture flux convergence, warm-rain processes, and rotating supercells all con-  
106 tributed to efficient precipitation processes in Harvey. Both *Wolff et al.* [2019] and *Brauer*  
107 *et al.* [2020] focused heavily on the precipitation processes and characteristics close to the  
108 Houston metropolitan area. In the current study, we focus on a broader region, including



109 the rainfall maximum near Beaumont, TX, with the aim of better understanding the vari-  
110 ability in rainfall in Harvey and the similarities and differences with rainfall in Florence.

111 The goal of this study is to characterize the nature of the extreme precipitation as-  
112 sociated with Harvey and Florence, taking advantage of polarimetric radar data and the  
113 dense rain gauge network (where available). We approach this study through two perspec-  
114 tives: 1) assessing the relative contributions of intense and long-lasting rainfall and 2) ex-  
115 amining characteristics of the near-surface rainfall inferred from polarimetric radar data.  
116 We first introduce our data processing methods (Sec. 2). Then we decompose the surface  
117 precipitation from Harvey into intense and light rain, examine how the rainfall unfolded  
118 at two representative rain gauge locations, and use the polarimetric radar data from Har-  
119 vey and Florence to compare the evolution and variability of the rainfall characteristics  
120 over a broad region of each storm (Sec. 3). We then use the polarimetric radar data to de-  
121 rive metrics to describe the bulk drop size distributions in each storm (Sec. 4). Finally,  
122 we propose hypotheses to explain the differences in the observed characteristics and the  
123 inferred microphysical processes from Harvey and Florence (Sec. 5).

## 124 **2 Data and Processing**

### 125 **2.1 Rain gauge data**

126 Hourly precipitation data come from the Automated Surface Observing System (ASOS)  
127 network. Due to strong winds and heavy rainfall, many rain gauges failed during Har-  
128 vey and Florence. We focus our analysis on the gauges that reported more than an av-  
129 erage of 20 hourly observations per day. This requirement allows us to include stations  
130 that are missing an occasional hourly observation, while excluding stations that are miss-  
131 ing too much data such that their statistics become meaningless. Most ASOS stations in  
132 Texas met the required number of hourly observations. Of the stations that did not meet  
133 the required number of hourly observations, all but one station failed for at least half the  
134 period, meaning the stations included were not sensitive to modest reductions in the re-  
135 quired number of hours. The vast majority of stations failed during Florence, particularly  
136 in North Carolina, where only two ASOS stations met our hourly observations require-  
137 ment. In South Carolina, roughly half the stations met our requirement and like Harvey,  
138 the stations that did not were also missing roughly half the observations. Unfortunately,

139 the lack of adequate stations from Florence prevents a complete comparison between rain  
140 gauge data from Florence and Harvey.

## 141 **2.2 Polarimetric radar data**

142 The radar data used in this study come from the operational polarimetric Next Gen-  
143 eration Weather Radar (NEXRAD) network. These radars have a wavelength of 10 cm  
144 (S band) and a  $1^\circ$  beamwidth, which allows for spatial resolution of order  $\sim 1$  km over  
145 large distances with minimal attenuation. The key benefit of polarimetric capabilities is  
146 the inference of bulk hydrometeor characteristics through the use of both horizontally and  
147 vertically polarized radio waves (*Bringi and Chandrasekar [2001]; Ryzhkov et al. [2005];*  
148 *Kumjian [2013]*). In Rayleigh scattering regimes, the radar reflectivity at horizontal polar-  
149 ization  $Z_H$  is proportional to the power returned from the backscatter due to both raindrop  
150 size and concentration, although it is more sensitive to the former due to the dependence  
151 on the sixth power of the drop diameter. The differential reflectivity  $Z_{DR}$  is defined as  
152 the difference between the reflectivity at horizontal and vertical polarization and is pro-  
153 portional to the median drop size of the DSD. The correlation coefficient  $\rho_{HV}$  is defined  
154 as the correlation of the horizontal and vertical pulses, with high values associated with  
155 nearly spherical targets like raindrops, and lower values typically associated with biolog-  
156 ical targets or mixed phase conditions within the pulse volume. The specific differential  
157 phase  $K_{DP}$  is the local change in the difference between the vertical and horizontal phase  
158 shifts and is proportional to the liquid water content, which depends on both the drop size  
159 and number concentration.

160 Located just west of Galveston Bay, the Houston radar (KHGX) captured the vast  
161 majority of Harvey's precipitation. Along the coast in central North Carolina, the More-  
162 head City radar (KMHX) was situated to the northeast of Florence's maximum precipi-  
163 tation, but unfortunately went offline after 1800 UTC on 15 September, missing the final  
164 24 h of rainfall. While the nearby Wilmington radar (KLTX) was better positioned than  
165 KMHX and operated over a longer period, it suffered from an apparent  $Z_{DR}$  bias and in-  
166 termittent dropouts (not shown). Although applying an estimated bias correction to the  
167 KLTX data produced results that were reasonably consistent with the KMHX radar, we fo-  
168 cus our analysis on the higher-quality KMHX data to provide a better comparison with the  
169 high-quality radar data from Harvey.

170 To ensure the best representation of near-surface characteristics, we restrict the ma-  
 171 jority of our analysis to data from the lowest elevation angle ( $0.5^\circ$ ) from each radar. We  
 172 retain the native polar coordinates of the radar to avoid interpolation, but the total and  
 173 fractional coverage of each polarimetric variable and derived metric category at each range  
 174 gate are weighted by the distance from the radar to account for beam spreading at larger  
 175 radii. We exclude data beyond 127 km from the radar to limit contamination by ice above  
 176 the melting layer. We note that by including all data within 127 km of KHGX our analy-  
 177 sis encompasses a larger area than the one used by *Wolff et al.* [2019], which focused on  
 178 the region immediately surrounding the Harris County Flood Warning System Network.

179 LROSE software (*Bell* [2019]) was first used to determine the most likely hydrome-  
 180 teor type according to the National Center for Atmospheric Research (NCAR) Particle ID  
 181 (PID) fuzzy logic algorithm (*Vivekanandan et al.* [1999]). We then isolated all radar pix-  
 182 els that were classified as light, moderate, or heavy rain and also had  $\rho_{HV}$  values between  
 183 0.95 and 1.0, inclusive, to retain only rain in our analysis. We performed an additional  
 184 subjective quality control by identifying persistent signals of reduced or elevated  $Z_{DR}$   
 185 along entire beams and removed those beams from the analysis. This processing resulted  
 186 in the removal of 20 beams (2.8%) from KHGX believed to have been compromised by  
 187 partial beam blockage, while no beams from KMHX were compromised.

188  $K_{DP}$  estimation in LROSE is based on an updated version of the *Hubbert and Bringi*  
 189 [1995] method. First,  $\phi_{DP}$  is unfolded to create a monotonically increasing field. The un-  
 190 folded  $\phi_{DP}$  field is smoothed using a finite impulse response (FIR) filter, similar to *Hub-*  
 191 *bert and Bringi* [1995], except using fewer iterations. Local bumps in the smoothed  $\phi_{DP}$   
 192 field due to backscatter differential phase ( $\delta$ ) are removed before another FIR is applied  
 193 for additional smoothing. Finally, the algorithm computes  $K_{DP}$  as the derivative of the  
 194 smoothed and corrected  $\phi_{DP}$ . While  $K_{DP}$  can be negative in ice regions, negative val-  
 195 ues are unphysical in rain and indicative of noise in the derivative calculation. To further  
 196 reduce the impact of noise on our analysis, we removed all radar gates where  $K_{DP}$  was  
 197 negative.

198 The *Cunningham et al.* [2013] technique was used to estimate  $Z_{DR}$  bias due to cal-  
 199 ibration errors by applying the technique to low-reflectivity regions in each radar sweep  
 200 and averaging together the individual biases over all sweeps where the number of samples  
 201 surpassed 500. With this technique, we calculated  $Z_{DR}$  bias corrections of -0.25 and -0.07

202 dB to KHGX and KMHX, respectively. At times, we also noticed odd jumps in the mean  
203  $Z_{DR}$  at KHGX, which were well outside the mean bias correction and we deemed to be  
204 unphysical. We removed 22 of the  $0.5^\circ$  sweeps that exhibited strong jumps in the mean  
205  $Z_{DR}$  compared to the surrounding times.

### 206 **2.3 Storm and shear data**

207 The storm intensity and location data for Harvey and Florence come from the Na-  
208 tional Hurricane Center Best Track dataset. The deep-layer wind shear magnitude and  
209 direction for Harvey come from the Statistical Hurricane Intensity Prediction Scheme  
210 (SHIPS) database (*DeMaria et al.* [2005]). Since the final predictors had not been released  
211 for 2018 at the time of this manuscript, we use the real-time SHIPS predictors for Flo-  
212 rence. In each case, the deep-layer shear is calculated from 850-200 hPa over radii be-  
213 tween 0 and 500 km from the storm center after the vortex has been removed.

## 214 **3 Rain gauge and polarimetric radar characteristics**

### 215 **3.1 Harvey rainfall**

224 Harvey first made landfall near Rockport, TX as a Category 4 hurricane (*Blake and*  
225 *Zelinsky* [2018]). Weak steering flow slowed and eventually reversed the inland movement,  
226 sending Harvey back to the Gulf of Mexico around 1200 UTC on 28 August. Although  
227 the *center* of Harvey remained at least 170 km from Houston, heavy rain accumulated  
228 near Houston and Beaumont between 0000 UTC on 25 August and 0000 UTC on 31 Au-  
229 gust as persistent onshore flow rose over a stationary front (*Blake and Zelinsky* [2018]).  
230 Overall, rain gauges near the coast measured the largest rainfall accumulations (Fig. 1a).  
231 Notably, Houston's George Bush Intercontinental Airport (KIAH) and Beaumont's Jack  
232 Brooks Regional Airport (KBPT) received 794 mm and 1207 mm in five days, respec-  
233 tively.

234 Event totals alone are insufficient for understanding extreme rainfall, which depends  
235 on both rain rate and event duration. Given Harvey's slow motion, it is clear that the long  
236 period of rainfall was an important factor. In an attempt to untangle the two effects, we  
237 isolate the frequency and fractional contribution of intense rain, which we define as a rain  
238 rate exceeding  $25.4 \text{ mm h}^{-1}$  ( $1 \text{ in h}^{-1}$ ). This threshold is consistent with the minimum rain  
239 rate of  $25 \text{ mm h}^{-1}$  used by *Hitchens et al.* [2013] to define heavy rain in their climatol-

240 ogy of hourly rain rates in the continental United States. *Hitchens et al.* [2013] show that  
241 such a rain rate is infrequent but not rare. Gauge-adjusted radar estimates of precipitation  
242 suggest  $10^6$  instances of  $25 \text{ mm h}^{-1}$  rain rates and  $10^2$ - $10^3$  instances of  $150 \text{ mm h}^{-1}$  rain  
243 rates occur each year, compared to over  $10^7$  instances of  $10 \text{ mm h}^{-1}$  rain rates (*Hitchens*  
244 *et al.* [2013]). For each station, we calculate the frequency of intense and light rain (ex-  
245 cluding hours with no rain) and their relative contributions to the total storm rainfall over  
246 the five-day period. By our definition, intense rain was infrequent and comprised no more  
247 than one-fourth of all raining hours, but the intense rain occurrence was closely associ-  
248 ated with the total rain (Fig. 1). Intense rain contributed nearly one-third of the total rain  
249 in the Houston metro area. Meanwhile, intense rainfall made up one-fourth of all hours  
250 with measurable rainfall at stations close to the Louisiana border and was responsible for  
251 nearly two-thirds of the total rain. The two stations near Beaumont, TX consistently saw  
252 a larger fraction of the total rainfall come from intense rain rates than all but one station,  
253 regardless of whether intense rainfall threshold was 20, 30, 40, or  $50 \text{ mm h}^{-1}$  (not shown).  
254 The only other station that had comparable fractional contribution of intense rainfall when  
255 the intense rainfall threshold was 40 or  $50 \text{ mm h}^{-1}$ , is located on the west side of Houston  
256 ( $29.62\text{N}$ ,  $-95.65\text{W}$ ). The gauge observed less rainfall than the Beaumont gauges and most  
257 of the Houston gauges and three hours of intense rain were able to make a larger frac-  
258 tional contribution. Both intense rain rates and the long duration were important, but their  
259 relative importance varied across Texas.

260 The spatial variability of rainfall evolution is highlighted in the precipitation time  
261 series for KIAH and KBPT (Fig. 1b, c). We focus on the 120-h (5-d) period after 0000  
262 UTC on 26 August as it encompasses most of the rain at each location (KIAH: 99%;  
263 KBPT: 97%). Overall, the hourly KIAH timeseries is consistent with the 15-minute ob-  
264 servations from the Harris County Flood Warning System gauges (*Wolff et al.* [2019]). To  
265 first order, Harvey's slow eastward movement dominates differences between the time se-  
266 ries in Houston and Beaumont. The daily rainfall at KIAH peaked on 27 August as the  
267 center of Harvey passed by Victoria, TX, placing the Houston metro area closer to the cir-  
268 culation center and directly in the onshore flow. Meanwhile, rainfall at KBPT peaked dur-  
269 ing the 12-h period surrounding 0000 UTC on 30 August as the center of Harvey moved  
270 southeast of Houston, when Beaumont was closer to the storm center. At both airports,  
271 the peak rainfall was strong and steady, but rain rates were more intense at KBPT. Before  
272 the peak rainfall, both locations observed intermittent intense rain suggesting embedded

273 convection within a broader region of stratiform precipitation. KBPT was in this regime  
274 for a longer period than KIAH and only light rain occurred at KIAH after Harvey pushed  
275 east. The period of light rain should not be disregarded, as it amounts to 20% of the total  
276 precipitation at KIAH. But the combination of prolonged convection and more intense rain  
277 rates led to more precipitation at KBPT than KIAH.

### 278 **3.2 Polarimetric characteristics of Harvey**

283 As discussed in the introduction, a key benefit of polarimetric radars is the ability  
284 to infer microphysical characteristics over a large area. A sample sweep from the KHGX  
285 radar shows widespread rain over southeast TX on 27 August (Fig. 2a). Radar reflectiv-  
286 ity values above 40 dBZ are frequent, but the echo intensities are not uniform and the  
287 NCAR PID algorithm identifies differences in the rain intensity (Fig. 2b). To examine  
288 the radar data through time, we select all radar gates classified by the PID as light, mod-  
289 erate, or heavy rain. In *Vivekanandan et al.* [1999], these categories roughly correspond to  
290 rain rates of  $< 10 \text{ mm h}^{-1}$ ,  $< 40 \text{ mm h}^{-1}$ , and  $> 40 \text{ mm h}^{-1}$ , respectively. These categories  
291 do not match our definitions of light and intense rainfall, since our primary usage of the  
292 PID categories is merely to identify likely raining radar echoes. Some infrequent echoes  
293 believed to be falsely identified as graupel are excluded from our analysis (Fig. 2b).

294 Once the raining areas are identified, we count the frequency of each PID rain cate-  
295 gory within 127 km of the radar. We weight the count for each radar gate by its distance  
296 from the radar to account for beam spreading at larger range distances and aggregate the  
297 weighted counts at each hour to create a time series of hourly PID counts (Fig. 2c). We  
298 include data over the ocean, but our results are not sensitive to the exclusion of offshore  
299 data points. Rain was widespread and peaked in coverage on 27 August. Harvey's exit  
300 from the radar domain is visible in the diminishing counts after 1200 UTC on 30 Au-  
301 gust. Occasional missing or bad radar sweeps yield gaps in the time series. To determine  
302 the relative importance of each PID category for each hour over the precipitating area,  
303 we examine the fractional coverage of the PID categories (Fig. 2d). Any periods when  
304 the weighted count failed to exceed  $2 \times 10^{10}$  were removed to ensure a sufficient sample  
305 size. Overall, light rain was the dominant PID category, but heavy and moderate rain were  
306 more frequent before 1200 UTC on 28 August. A slight resurgence in moderate rain oc-  
307 curred between 0000-0600 UTC on 30 August.

Using these rain categories to identify likely precipitating echo, we examine the distributions of  $Z_H$ ,  $Z_{DR}$ , and  $K_{DP}$  to determine the dominant microphysical characteristics. One prominent feature of the distributions is the modest values of each polarimetric variable, consistent with a midsize-drop dominated regime (Fig. 3).  $Z_H$  and  $Z_{DR}$  values lie mostly below 45 dBZ and 2 dB, respectively, consistent with prior studies (*Brown et al. [2016]*, *Wang et al. [2016]*, *Wolff et al. [2019]*, and *Brauer et al. [2020]*).  $K_{DP}$  values infrequently surpass the  $0.3^\circ \text{ km}^{-1}$  threshold often used for  $K_{DP}$ -based rain rate algorithms, similar to  $K_{DP}$  values estimated from disdrometer observations in tropical convection over the Indo-Pacific warm pool (*Thompson et al. [2018]*). Despite modest values overall,  $Z_H$ ,  $Z_{DR}$ , and  $K_{DP}$  values were larger in the first half of the event, consistent with the increased prevalence of moderate and heavy rain identified by the PID algorithm. The  $K_{DP}$  values in our analysis cover the same range as the values obtained by *Wolff et al. [2019]*, but our distribution skews towards weaker values. Our area of analysis is roughly a factor of 4 larger than *Wolff et al. [2019]* that focused on a rectangle located in the northwest quadrant of the radar domain to compare radar-estimated rain rates with rain gauges in the Harris County Flood Warning System Network. We also use a slightly different algorithm to estimate  $K_{DP}$ . Since the estimated  $K_{DP}$  values span a similar range of values, we hypothesize that the different regions of analysis contribute the most to the difference in  $K_{DP}$  values.

The greatest shift to larger reflectivity values in the full-domain distributions occurred in two periods: 0600 UTC on 27 August through 1200 UTC on 28 August and 0000-0600 UTC on 30 August (Fig. 3a). Despite similar reflectivity distributions,  $Z_{DR}$  values were greater in between 0300-0900 UTC 27 August, indicating larger drops were responsible for the higher reflectivities (Fig. 3b). The simultaneous increase in  $Z_H$  and  $Z_{DR}$  values on 27 August were noted by both *Wolff et al. [2019]* and *Brauer et al. [2020]*. Although the larger-domain hourly timeseries show that the most intense values only persisted for half the day. In contrast to the high  $Z_{DR}$  values on 27 August, the high reflectivities on 30 August coincided with reduced  $Z_{DR}$  values, indicative of smaller drops. In between was a 36-h period (1200 UTC on 28 August-0000 UTC 30 August) of weak reflectivity values across the domain. The coincident decrease in  $Z_{DR}$  and  $K_{DP}$  values around 1200 UTC 28 August suggests a decrease in the drop size. There was a similar weakening of the polarimetric variables after the second heavy rain peak (~1200 UTC 30 August), but the echo coverage at this time was minimal (Fig. 2c).

349 To get a sense of how the polarimetric data relate to the rain gauge data at key time  
 350 periods, we compare radar data surrounding KIAH and KBPT during hours when the sta-  
 351 tions experienced their heaviest rainfall. Figure 4 shows snapshots of gridded radar data  
 352 near KIAH and KBPT during those hours of intense rainfall. Radar data was gridded with  
 353 LROSE Radx2Grid software with horizontal grid spacing of 1 km and vertical grid spac-  
 354 ing of 0.5 km below 3 km altitude and 1.0 km above 3 km altitude. Near KIAH, the 0421  
 355 UTC radar observation on 27 August coincided with an hourly rain gauge observation of  
 356 nearly 50 mm. Reflectivity values at 1 km altitude within 32 km of the gauge ranged from  
 357 25-50 dBZ, while those immediately next to the gauge exhibited a narrower range between  
 358 35-45 dBZ (Fig. 4a). Meanwhile, near KBPT, the 0431 UTC radar observation on 30 Au-  
 359 gust coincided with an hourly rain gauge observation of almost 100 mm, which was the  
 360 highest rain rate at that location. Although 1 km altitude reflectivity values surrounding  
 361 the gauge spanned a similar range as observed near KIAH days prior, reflectivity values  
 362 were more homogeneous, with a large area of reflectivity values exceeding 40 dBZ (Fig.  
 363 4b). Comparing the vertical profiles of  $Z_H$ ,  $Z_{DR}$ , and  $K_{DP}$  shows that each variable was  
 364 generally more intense near KBPT throughout the atmosphere, with the exception of  $Z_{DR}$   
 365 values above 6 km altitude, where the spread was large (Fig. 4c-e). All three polarimet-  
 366 ric profiles increase towards the surface near KBPT, consistent with enhanced collision-  
 367 coalescence processes (Kumjian and Prat [2014]). Meanwhile,  $Z_H$  increases toward the  
 368 surface near KIAH, but the  $K_{DP}$  increase toward the surface is weaker and  $Z_{DR}$  exhibits a  
 369 lot of spread, including both positive and negative slopes. The variability in vertical struc-  
 370 ture is not surprising given the heterogeneous low-level reflectivity field near KIAH at  
 371 this time; Brauer *et al.* [2020] also showed that vertical polarimetric profiles over nearby  
 372 downtown Houston varied substantially on 27 August.

### 373 3.3 Florence rainfall

377 The outer rainbands of Florence approached North Carolina on 13 September and  
 378 Florence made landfall in southeast North Carolina on 14 September as a Category 1  
 379 storm. Similar to Harvey, forward motion slowed as Florence crept slowly into South Car-  
 380 olina before accelerating northward late on 16 September (Fig. 5). Unfortunately, exten-  
 381 sive rain gauge outages in the region of maximum rainfall preclude a breakdown of rain  
 382 rates for Florence, as indicated by the number of stations with too many hours of missing  
 383 data in Fig. 5. The available data reveal two big differences from Harvey, however. First,



384 Florence was a shorter event as most gauges observed rainfall for approximately three  
 385 days. Second, the available rain rates from Florence are comparable, yet slightly weaker  
 386 than the rain rates in Fig. 1b,c, although rain rates likely strengthened on 15 September  
 387 when most gauges were offline. As discussed in section 2, the Morehead City, NC radar  
 388 (KMHX) also went offline midway through the event (~1800 UTC 15 September), missing  
 389 the final 12-24 hours of rainfall over North Carolina. Despite missing data, we think the  
 390 high radar data quality, sufficient length of the data record (> 48 h), and similar statistics  
 391 to bias-corrected KLTX radar data allow for a reasonable comparison of Florence's rainfall  
 392 characteristics derived from KMHX and the rainfall characteristics from Harvey.

### 393 **3.4 Polarimetric characteristics of Florence**

395 Similar to Harvey, polarimetric data from Florence are dominated spatially and tem-  
 396 porally by the light rain PID category (Fig. 6a). The magnitude of range-weighted counts  
 397 is similar, although KMHX observed more offshore pixels than KHGX. Similar to Harvey,  
 398 our results are not sensitive to the exclusion of offshore data points. The range-weighted  
 399 counts increased slowly on 13 September as Florence approached the east coast, before  
 400 plateauing on 14 September (Fig. 6b). As Florence moved southwest away from the radar,  
 401 the range-weighted counts decreased on 15 September. At the same time, the fractional  
 402 coverage of moderate and heavy rain categories increased (Fig. 6c).

404 The polarimetric distributions exhibit similar tropical cyclone characteristics to Har-  
 405 vey, although most values are generally lower. Reflectivity values seldom exceed 45 dBZ,  
 406 differential reflectivity values never exceed 2 dB, and  $K_{DP}$  values are lower than those  
 407 seen during Harvey (Fig. 7). The lower polarimetric values are consistent with the values  
 408 observed in an inner rainband in *Wu et al.* [2018]. Greater coverage of the moderate inten-  
 409 sity values corresponds to reduced coverage of heavier rain rates in comparison to Harvey.  
 410 Additionally, the polarimetric distributions from Florence gradually shift to higher val-  
 411 ues after 0600 UTC on 15 September in contrast to the more episodic nature of Harvey's  
 412 variability. The increase in coverage of  $Z_H$ ,  $Z_{DR}$ , and  $K_{DP}$  above 40 dBZ, 1 dB, and 0.3  
 413 ( $^{\circ}\text{km}^{-1}$ ), respectively, are the closest the polarimetric values from Florence approach the  
 414 statistics from the first 60 hours of Harvey (Fig. 3). Even if these statistics persisted for  
 415 the final 24 hours that Florence impacted North Carolina, the length of influence by heav-  
 416 ier rain rates would only amount to 48 hours. The amount of intense polarimetric values  
 417 in Florence was reduced spatially and temporally in comparison to Harvey.

424 Although most rain gauges failed in Florence, we can still examine representative  
 425 vertical profiles of polarimetric quantities. Due to the lack of suitable rain gauge com-  
 426 parisons, we chose instead to highlight the vertical structure of two different TC features:  
 427 a convective outer rainband and an inner rainband. Figure 8 shows snapshots of the hor-  
 428 izontal and vertical structure of two rainbands occurring simultaneously in Florence on  
 429 14 September 2018. The outer rainband exhibits a similar structure as the examples from  
 430 Harvey in Fig. 4, with heterogeneous horizontal patterns of reflectivity at 1 km and re-  
 431 flectivity peaking between 45-52 dBZ. The vertical profiles of the polarimetric variables  
 432 in the outer rainband shown in Fig. 8c-e are similar to those near KBPT during Harvey  
 433 on 30 August, where  $Z_H$ ,  $Z_{DR}$ , and  $K_{DP}$  all increase toward the ground as collision-  
 434 coalescence processes are enhanced (*Kumjian and Prat* [2014]). Meanwhile, the inner  
 435 rainband shows more modest increases in  $Z_H$  and  $K_{DP}$  near the surface, and the  $Z_{DR}$   
 436 profile varies little below 4 km. These localized profiles combined with the examples from  
 437 Harvey underscore the variety of microphysical processes that can exist in a TC, consis-  
 438 tent with prior studies (*Wang et al.* [2016]; *Didlake and Kumjian* [2018]; *Wu et al.* [2018]).  
 439 A more systematic examination of the vertical polarimetric profiles with respect to both  
 440 TC features and surface rainfall is beyond the scope of this study but is recommended for  
 441 future work.

#### 442 **4 Estimated DSD characteristics**

443 To better understand the covariability of the rain drop size and number concentra-  
 444 tion, we retrieve specific metrics related to the DSD from the radar data. A DSD can be  
 445 approximated as a gamma distribution with the following form:

$$446 \quad N(D) = N_0 D^\mu \exp(-\Lambda D) \quad (1)$$

447 where  $N_0$  is the intercept parameter,  $D$  is the drop diameter,  $\mu$  is the shape parameter, and  
 448  $\Lambda$  is the slope parameter (*Ulbrich* [1983]). The distribution can be normalized (*Willis*  
 449 [1984]) to compare different DSDs using a modified intercept parameter proportional to  
 450 the liquid water content and median drop diameter:

$$451 \quad N_W = \frac{1.81 \times 10^5 LWC}{\pi \rho_L D_0^4} \quad (2)$$

452 where  $N_W$  ( $\text{m}^{-3} \text{mm}^{-1}$ ) is the normalized intercept parameter,  $LWC$  ( $\text{g m}^{-3}$ ) is the liquid  
 453 water content,  $\rho_L$  ( $\text{g cm}^{-3}$ ) is the density of liquid water, and  $D_0$  (mm) is the median vol-  
 454 ume diameter.  $N_W$  and  $D_0$  can be considered proxies for the number of drops and the

455 median drop size. Although an assumed gamma distribution cannot represent all observed  
 456 DSDs, gamma distributions are prevalent within the literature and are an effective way of  
 457 identifying meaningful spatial and temporal changes in the type of DSD (*Willis* [1984];  
 458 *Bringi et al.* [2003]; *Chang et al.* [2009]; *Thompson et al.* [2015]; *Wang et al.* [2016]; *Za-*  
 459 *grodnik et al.* [2018]). Without disdrometer measurements of  $D_0$  and  $N_W$ , we use the  
 460 CSU RadarTools package (*Lang et al.* [2019]) to estimate these parameters from polari-  
 461 metric data following the algorithm described by *Bringi et al.* [2015] where  $D_0$  and  $N_W$   
 462 are related to  $Z_{DR}$  and  $Z_H$  through the following equations:

$$463 \quad D_0 = \begin{cases} 0.0536Z_{DR}^3 - 0.1971Z_{DR}^2 + 0.6261Z_{DR} + 1.0815, & Z_{DR} \geq 1\text{dB} \\ 0.0424Z_{DR}^4 - 0.4571Z_{DR}^3 + 0.6215Z_{DR}^2 + 0.457Z_{DR} + 0.8808, & Z_{DR} < 1\text{dB} \end{cases} \quad (3)$$

$$464 \quad N_W = 19.76 \frac{Z_H}{D_0^{7.46}} \quad (4)$$

465 We exclude data where  $Z_{DR}$  is below -0.5 dB, though such data points are infrequent.  
 466 To understand how the DSDs are broadly related to water content, we obtain a theoret-  
 467 ical estimate of LWC by rearranging equation 2. Infrequent big drop and numerous small  
 468 drop DSDs can produce similar LWC values, despite distinct radar signatures and forma-  
 469 tion processes (Fig. 9). As discussed by *Chang et al.* [2009], precipitation radar estimates  
 470 such as those from S-band can miss the numerous ( $\log_{10}(N_W) > 4$ ), small drop ( $D_0 < 1$   
 471 mm) DSDs due to limitations in radar sensitivity, which can cause an underestimate of the  
 472 LWC from precipitation radar estimates. Despite these uncertainties, by using the same  
 473 algorithm and radar wavelength to estimate  $N_W$  and  $D_0$  in Harvey and Florence, we can  
 474 quantitatively compare the drop size characteristics of the two TCs, while qualitatively  
 475 comparing with previous studies that use disdrometer measurements or different radar re-  
 476 trieval techniques.  
 477

486 Figure 9a shows that numerous small drop DSDs were common in Harvey, which  
 487 is broadly consistent with previous studies on TC DSDs (*Tokay et al.* [2008]; *Chang et al.*  
 488 [2009]; *Wolff et al.* [2019]). The retrieval suggests Harvey's dominant DSD type lies in  
 489 between small drop DSDs from maritime convection (*Bringi et al.* [2003]; *Thompson et al.*  
 490 [2015]) and large drop DSDs that can occur in continental, wintertime precipitation (*Za-*  
 491 *grodnik et al.* [2018]), but the spread differs from other polarimetric estimates of DSD  
 492 parameters in TCs. The retrieved DSDs exhibit lower concentrations than *Wang et al.*  
 493 [2016] by an order of magnitude ( $\Delta 1$  in  $\log_{10}(N_W)$  space) and the median drop diameters  
 494 are approximately 0.5 mm smaller than the values retrieved by *Chang et al.* [2009]. That

495 being said, we emphasize that these two studies use a different technique (*Zhang et al.*  
 496 [2001]) that estimates  $N_0$  and  $\Lambda$  in Eq. 1 rather than the parameters estimated here in Eq.  
 497 2. To test the sensitivity to the retrieval algorithm, we calculated the DSD parameters us-  
 498 ing the relationship obtained by *Brandes et al.* [2004]. The resulting  $N_W$ - $D_0$  distributions,  
 499 the temporal evolution of each storm, and the differences between Harvey and Florence  
 500 were quite similar using either the *Bringi et al.* [2015] or *Brandes et al.* [2004] retrievals  
 501 (not shown). As a result, we feel confident that our comparisons between Harvey and Flo-  
 502 rence using the same retrieval technique reveal differences in rainfall characteristics and  
 503 microphysical processes between the two storms. Further comparisons with other storms  
 504 may reflect retrieval technique uncertainty in addition to differences in microphysical pro-  
 505 cesses.

506 The derived DSD parameters in Florence span similar values to Harvey, and show a  
 507 qualitatively similar joint probability distribution (Fig. 9b). Some distinct differences are  
 508 apparent however, with a reduction in the spread of the observed DSD variability and a  
 509 shift towards smaller drops (Fig. 9c). The most common DSDs have an estimated  $D_0$  at  
 510 or below 1.0 mm, with relatively high number concentration. The overall similarity be-  
 511 tween the Harvey and Florence probability distributions suggest comparable microphysical  
 512 processes in general in both events but with more stronger convection in Harvey.

513 To provide more insight into the variability of the observed DSDs, we divide the  
 514 joint  $N_W$ - $D_0$  distribution into four quadrants using boundaries of  $3.5 \text{ m}^{-3} \text{ mm}^{-1}$  and 1.2  
 515 mm (Fig. 9). These thresholds are arbitrary, but approximate the midpoint of each distri-  
 516 bution. Not only does each quadrant have a different combination of median drop size and  
 517 number concentration, but the thresholds separate the distributions by theoretical LWC val-  
 518 ues. In particular, the difference in the theoretical LWC between the lower-left (low LWC)  
 519 and upper-right (high LWC) quadrants is substantial. The following quadrant descriptions  
 520 are defined for the purposes of comparison: low-concentration small drops (quadrant SL),  
 521 low-concentration medium drops (quadrant ML), high-concentration small drops (quadrant  
 522 SH), and high-concentration medium drops (quadrant MH).

525 In Harvey, all four DSD quadrants coexisted over the radar domain at all times but  
 526 their relative coverage varied in time (Fig. 10). The fractional coverage of larger  $D_0$  val-  
 527 ues achieved temporary maxima midday on 26 and 27 August and temporary minima mid-  
 528 day on 28 August (Fig. 10a).  $N_W$  variability over time was smaller than the variability

529 in  $D_0$  (Fig. 10b). Although variable in time, the ML and SH quadrants are slightly more  
 530 frequent throughout Harvey, while the SH quadrant is the most common. (Fig. 10c).

531 Overall, two notable regime shifts in the DSDs occurred during Harvey's evolution.  
 532 First, a shift to larger drops occurred on 27 August, characterized by a maximum in  $D_0$   
 533 values, increased quadrant MH coverage, and a maximum in quadrant ML coverage (Fig.  
 534 10). At the same time, there is a shift to lower  $N_W$  values. These characteristics are con-  
 535 sistent with the conclusions of *Wolff et al.* [2019] and *Brauer et al.* [2020]. This regime  
 536 was short-lived, lasting only from 0400-01700 UTC on 27 August. The surge in medium  
 537 drop coverage is similar to observations of a convective cell in an outer rainband in Ty-  
 538 phoon Nida (2016) where bigger drops were prominent (*Wu et al.* [2018]). In *Wu et al.*  
 539 [2018], increased  $Z_H$ ,  $Z_{DR}$ , and  $K_{DP}$  values were found beneath strong, deep ascent and  
 540 high concentrations of graupel and aggregates, which indicated the importance of melt-  
 541 ing ice aloft to bigger raindrops below. Due to differences in methodology, it is unclear  
 542 whether the drops in Nida (2016) would fall into our ML or MH quadrants. In the study  
 543 by *Wang et al.* [2016] of Typhoon Matmo (2014), they did not observe a noticeable in-  
 544 crease in low-concentration medium drops (our ML quadrant) but did report a shift to-  
 545 wards higher number concentration (our MH quadrant) within a convective rainband. We  
 546 note that the increase in medium drops observed by KHGX radar occurred around the  
 547 same time that KIAH airport rain gauge was receiving its peak rain rates and the radar  
 548 gates closest to the gauge showed a greater occurrence of the high-concentration drop  
 549 quadrants (not shown). The fine scale spatial and temporal variability confirms the chal-  
 550 lenge of TC precipitation forecasts to accurately predict local maxima embedded within a  
 551 broad envelope of long duration light to moderate rain.

552 The second regime shift exhibited greater coverage by the high-concentration DSDs  
 553 (SH and MH) during two periods: 1800 UTC on 27 August through 1200 UTC on 28  
 554 August and 0000-0600 UTC 30 August. In particular, the numerous, small drop DSDs  
 555 (the SH quadrant) reached peak coverage. This smaller drop regime was noted by *Wolff*  
 556 *et al.* [2019] on 28 August, although they did not analyze data from 30 August as the rain  
 557 had exited Harris County. The prevalence of higher concentration of small to medium  
 558 size drops is likely associated with an enhanced warm rain process. The extended pe-  
 559 riod that rain was within range of the radar is a likely contributor to the large variety of  
 560 near-surface microphysical characteristics compared to the *Wang et al.* [2016] and *Wu et al.*  
 561 [2018] studies (4.5 days in the current study compared to less than 12 hours in the latter).

562 The lengthy duration enabled multiple regions of the TC with different precipitation char-  
563 acteristics to pass over the radar.

564 The derived DSDs in Florence evolved more gradually than in Harvey, with reduced  
565 temporal variability but an increase in the medium drops (the ML and MH quadrants) that  
566 became more prominent after 0600 UTC on 15 September. The high number concentra-  
567 tion DSDs (the SH and MH quadrants) have the greatest coverage, but the MH quadrant  
568 is less frequent than was seen for Harvey. The reduced frequency of the MH quadrant and  
569 increased frequency in the SH quadrant is due to generally weaker  $Z_H$  and  $Z_{DR}$  (cf. Fig.  
570 7a,b) and indicates a general shift to smaller drop sizes over time.

571 Comparing the available data from Harvey and Florence reveals broad commonali-  
572 ties in microphysics of TC rainfall but important differences in the event length and local  
573 rain intensity. Harvey lingered over Texas for a longer amount of time and radar data indi-  
574 cate that the heaviest rain rates were more frequent and long-lasting than in Florence. Po-  
575 larimetric data suggest both storms had a prevalence of small-to-medium sized raindrops  
576 ( $\leq 2$  mm) in moderate to high concentrations, but there was more temporal variability of  
577 the DSDs in Harvey than in Florence. We next examine some of the environmental and  
578 structural features that may have led to these differences.

## 581 **5 Discussion**

585 We hypothesize that several key factors were influential in producing the differences  
586 in observed rainfall in Harvey and Florence. First, Harvey and Florence had different in-  
587 tensities during their prolonged rain events. Although intensity does not directly correlate  
588 with storm structure, Harvey weakened more rapidly than Florence and spent more time as  
589 a tropical storm (Fig. 12a). At the same time, the vertical wind shear surrounding Harvey  
590 was stronger than for Florence, which is known to produce azimuthal variations in precip-  
591 itation. Deep-layer shear impacts vertical motion, and thus rainfall, in part by tilting the  
592 vortex. Vortex tilt induces a wavenumber-1 asymmetry of potential temperature, leading  
593 to persistent anomalies of isentropic ascent and descent (*Jones [1995]*). In addition, deep-  
594 layer shear imposes a wavenumber-1 pattern of radial flow and vorticity advection, which  
595 is roughly balanced by vortex compression or stretching (*Bender [1997]*).

599 As a weak storm under strong shear, Harvey was more asymmetric and disorganized.  
600 Initially, the strongest precipitation was found within convective outer rainbands that im-

601 pacted the Houston metropolitan region (Fig. 13a). The intense rainbands in Harvey share  
602 some similarities with the typhoon rainbands studied by *Wang et al.* [2016] and *Wu et al.*  
603 [2018] that indicated a prevalence of larger drops associated with stronger, deeper convective  
604 precipitation associated with 'outer' rainbands (*Yu and Tsai* [2013]; *Tang et al.* [2014];  
605 *Tang et al.* [2018]). One day later, hours after the center of Harvey reentered the Gulf of  
606 Mexico, precipitation weakened and became more scattered (Fig. 13b). As Harvey moved  
607 northeast, the precipitation strengthened and became more uniform (Fig. 13c). Finally,  
608 as the center of Harvey moved into Louisiana, the echo area was restricted to the region  
609 surrounding Beaumont, TX; the reduced echo area and higher reflectivity (Fig. 13d) were  
610 associated with the resurgence of the heavy rain PID category and stronger polarimetric  
611 variables on 30 August (cf. Figs. 2 and 3). Another notable feature in Harvey was the  
612 relative lack of precipitation offshore. This structure suggests a strong influence of both  
613 onshore flow and persistent southwesterly shear. Under this shear orientation, southeast  
614 Texas spent a great deal of time in the downshear quadrants of Harvey, which are often  
615 characterized by enhanced overall rainfall, convective precipitation near the eyewall, and  
616 a transition from convective to stratiform precipitation in outer rainbands (*Corbosiero and*  
617 *Molinari* [2002]; *DeHart et al.* [2014]; *Hence and Houze* [2011]; *Hence and Houze* [2012];  
618 *Reasor et al.* [2013]).

622 In contrast to Harvey's more asymmetric structure, we hypothesize that a slower decay  
623 rate and weaker shear favored a more resilient and axisymmetric structure for Florence.  
624 Increased rainfall axisymmetry is expected for stronger storms experiencing weaker  
625 deep-layer shear (*Chen et al.* [2006]). The eyewall and rainbands remained well-defined as  
626 Florence pushed through the KMHX domain (Fig. 14). The heaviest rain in Florence was  
627 mostly restricted to the eyewall and inner rainbands, with larger areas of weaker reflectivity  
628 values and non-convective precipitation similar to the inner rainband in Fig. 8b. Unlike  
629 Harvey, heavy rainfall in Florence did not preferentially occur onshore. Although hints  
630 of heavier rainfall onshore exist in the eyewall and outer rainband on 14 September (Fig.  
631 14b), a band of heavier rainfall is present offshore on 15 September (Fig. 14c). Determining  
632 which structures are due to processes associated with the coastal effects is challenging  
633 without numerical simulations, and should be a focus of further study.

634 During Florence's landfall, the shear direction veered from southerly to westerly  
635 through 0000 UTC on 14 September (Fig. 12c), but the shear magnitude remained steady  
636 around 10 kts through 1800 UTC on 14 September. The concurrent axisymmetric struc-

637 ture suggests that Florence was more resilient to the influence of deep-layer shear. The  
638 relative axisymmetry of Florence during the slow decay was reflected in the persistent  
639 fractions of DSD types. The slow increase of large-sized DSDs coincided with Florence's  
640 departure from the radar domain as a solitary rainband dominated the remaining precipi-  
641 tating area (Fig. 14c). Since different TC features exhibit different types of DSDs (*Yu and*  
642 *Tsai* [2013]; *Tang et al.* [2014]; *Wang et al.* [2016]; *Tang et al.* [2018]; *Wu et al.* [2018]), a  
643 gradual transition from a regime that includes eyewall, inner rainband, and outer rainband  
644 rainfall (e.g., 14 September in Fig. 14b) to predominantly outer rainband rainfall likely  
645 contributes to the changing DSD type fractions. Overall, the polarimetric data suggest that  
646 strong convection was less widespread in Florence than in Harvey. The stronger TC in-  
647 tensity and resilience to deep-layer shear are believed to limit the intense convection to  
648 smaller areas in the eyewall and rainbands compared to the stronger vertical motion forc-  
649 ing downshear in Harvey.

650 Vertical wind shear and intensity are not the only possible factors that could ex-  
651 plain the difference in precipitation between the two storms. Environmental conditions  
652 such as dry air, thermodynamic stability, and sea surface temperatures also affect cloud  
653 morphology and the resulting precipitation. Offshore sea surface temperatures in the Gulf  
654 of Mexico prior to Harvey were approximately 1°C warmer than offshore temperatures  
655 for Florence (not shown) and could have increased low-level moisture in Harvey. While  
656 these factors may also play a role, the contrasting storm structures and consistency with  
657 expected patterns of precipitation due to vertical wind shear suggest that the combined  
658 effects of shear and intensity played an important role in the rainfall differences. Fur-  
659 ther analysis with high-resolution thermodynamic observations and numerical simulations  
660 would be required to evaluate their relative contributions and are beyond the scope of this  
661 study.

## 662 **6 Conclusions**

663 In this study, rain gauge and polarimetric radar data were analyzed to better under-  
664 stand the characteristics of the record-setting rainfall from Hurricanes Harvey (2017) and  
665 Florence (2018) and to identify microphysical similarities and differences between these  
666 two storms. From our analysis, we draw the following conclusions:



- 667 1. Polarimetric data indicated that the microphysical characteristics of both storms  
668 were not drastically different from the observations in previous tropical cyclones.  
669 On average, both Harvey and Florence exhibited ‘typical’ tropical cyclone DSDs  
670 with high concentrations of small-to-medium sized raindrops.
- 671 2. High-concentration, midsize DSDs were more common in Harvey and indicate  
672 stronger rain rates over a larger area for a longer time period. The dominant DSD  
673 type (e.g., large  $N_W$  or  $D_0$  DSDs) in Harvey showed substantial variability over  
674 time. In contrast, inferred DSDs from Florence generally had smaller median diam-  
675 eters than the DSDs from Harvey and had less variability.
- 676 3. Both the long event duration and strong rain rates were contributing factors to the  
677 record-breaking rainfall in Harvey, but their relative importance varied spatially.  
678 The slowly moving center determined the amount of time any one location spent  
679 in onshore flow and was one of the primary contributors to the frequency of heavy  
680 rain rates and the record-breaking total accumulated rainfall. In contrast, Florence’s  
681 more direct track and shorter duration did not produce as distinct a difference in on  
682 and offshore precipitation or the spatial or temporal variability in intense rain rates.
- 683 4. Stronger vertical wind shear and a more rapid decay in intensity contributed to  
684 a more asymmetric and disorganized Harvey, whereas weaker wind shear and a  
685 slower decay in intensity contributed to a more resilient and axisymmetric Florence.  
686 These factors contribute to overall weaker convection in Florence associated with  
687 ‘inner’ rainbands, while Harvey’s stronger convection more resembled ‘outer’ rain-  
688 bands.

689 It is well-known that slow-moving tropical cyclones are capable of producing ex-  
690 treme rainfall, but this study further underscores the complexity inherent in tropical cy-  
691 clone rainfall. Rain rates and the dominant microphysical processes vary sharply over  
692 short distances, depend on the details of the TC structure, and are influenced by the en-  
693 vironment. The changing DSDs in Harvey and Florence suggests changes in the relative  
694 contributions of rain and ice phase processes, consistent with previous TC studies (*Wang*  
695 *et al.* [2016]; *Didlake and Kumjian* [2017]; *Kalina et al.* [2017]; *Wu et al.* [2018]). Abun-  
696 dant small drops suggest the dominance of warm rain processes, while the presence of  
697 larger drops suggest more vigorous or deeper convection with more contributions from  
698 melting graupel to the overall DSD (*Wu et al.* [2018]). The current results suggest that  
699 multiple microphysical pathways to heavy rainfall exist within the same TC, and additional

700 research is needed to determine which pathways will dominate in different regions and  
701 landfall scenarios.

702 Understanding the key factors in specific extreme rainfall events can help identify  
703 commonalities and differences in future extreme events. Recent studies suggest that trop-  
704 ical cyclone rainfall is expected to increase in a warmer climate, and that climate change  
705 may have influenced Harvey's rainfall through increased water vapor, higher ocean heat  
706 content, or slower storm motion (*Emanuel [2017]; van Oldenborgh et al. [2017]; Risser*  
707 *and Wehner [2017]; Trenberth et al. [2018]*). While the role of climate change in the two  
708 events presented here is outside the scope of this study, documenting the characteristics  
709 of their rainfall is crucial to understanding how such characteristics might change in the  
710 future. In particular, the relative spatial and temporal contributions of various microphysi-  
711 cal processes to the total rainfall are not fully understood. A specific process might be an  
712 efficient producer of strong precipitation, but may have a limited impact in the context of  
713 the full event and area of impact. Concurrent surface, in situ, and remote observations of  
714 microphysical processes will improve our understanding of and ability to forecast tropical  
715 cyclone rainfall in the current and future climate.

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722 maintained by the Iowa Environmental Mesonet (<https://mesonet.agron.iastate.edu/request/download.phtml>). The LROSE software is available for download online (<http://lrose.net>).

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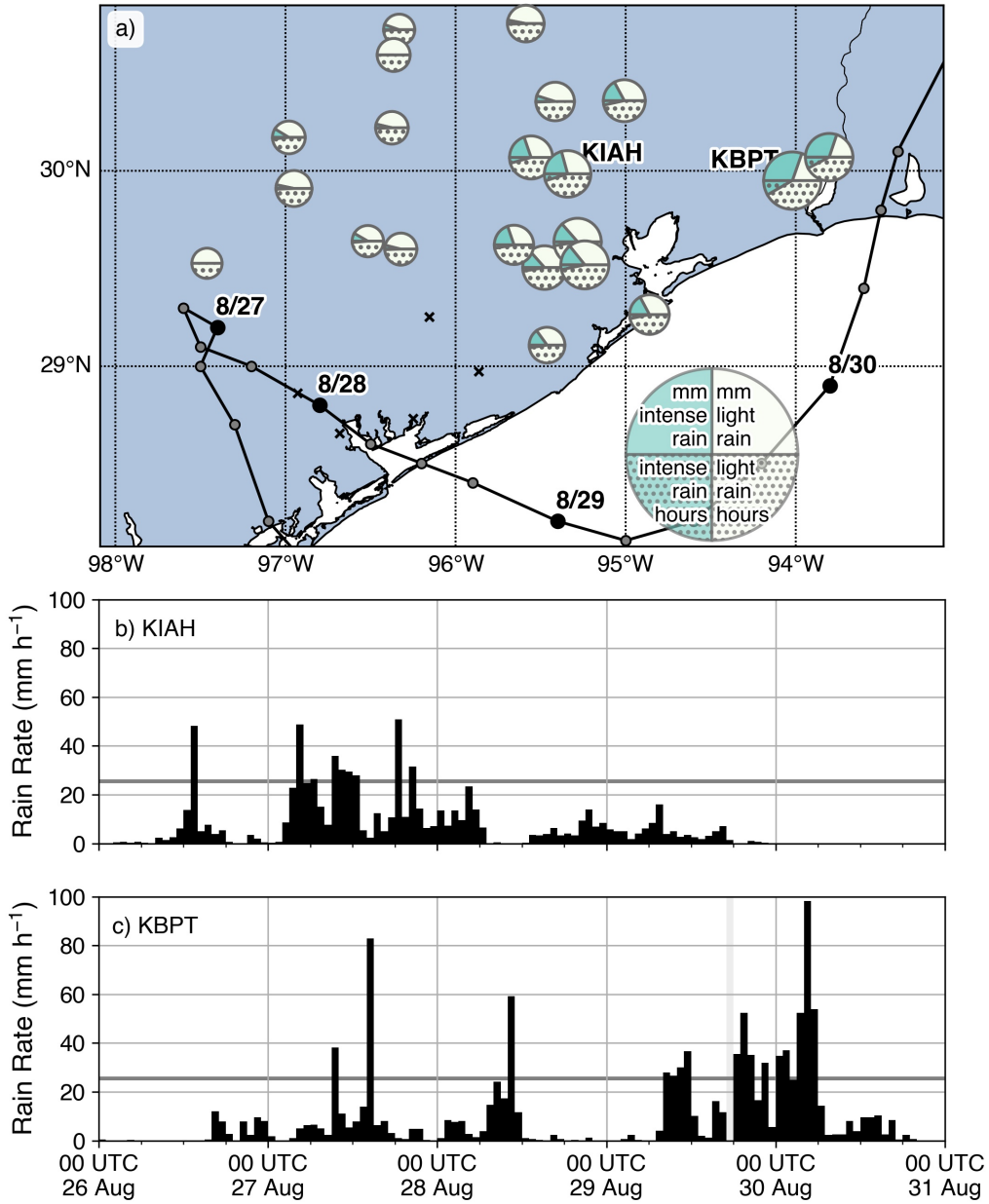
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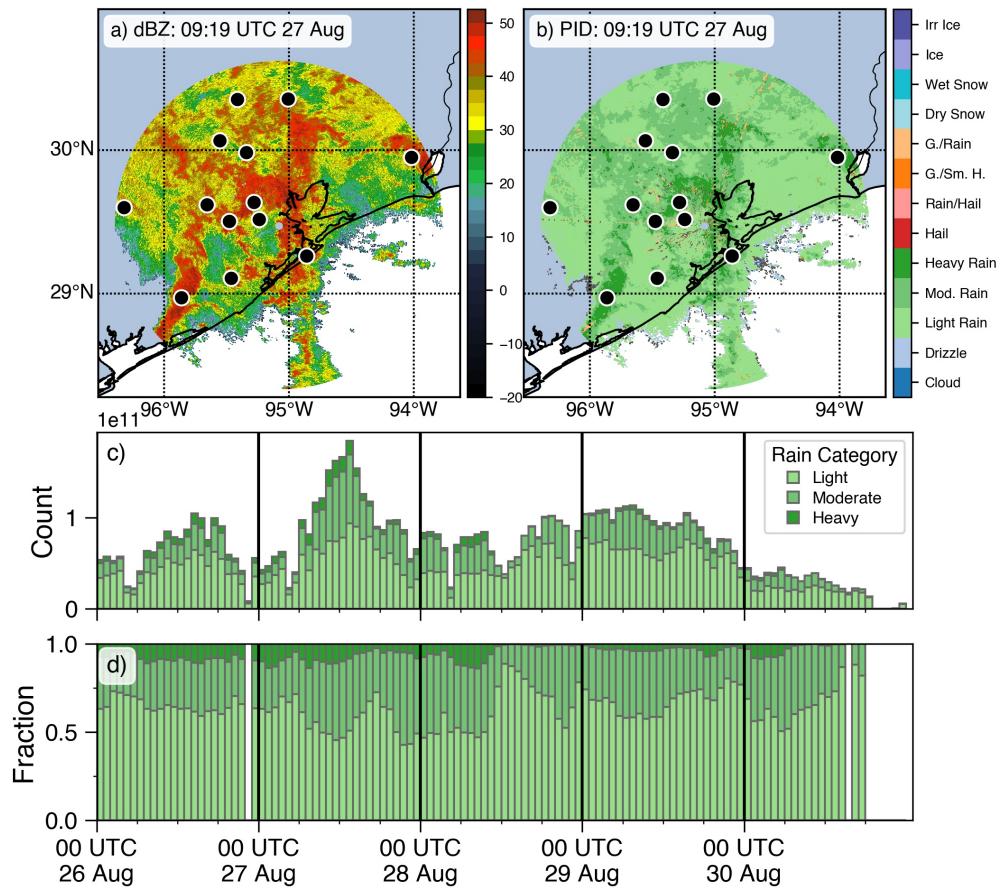
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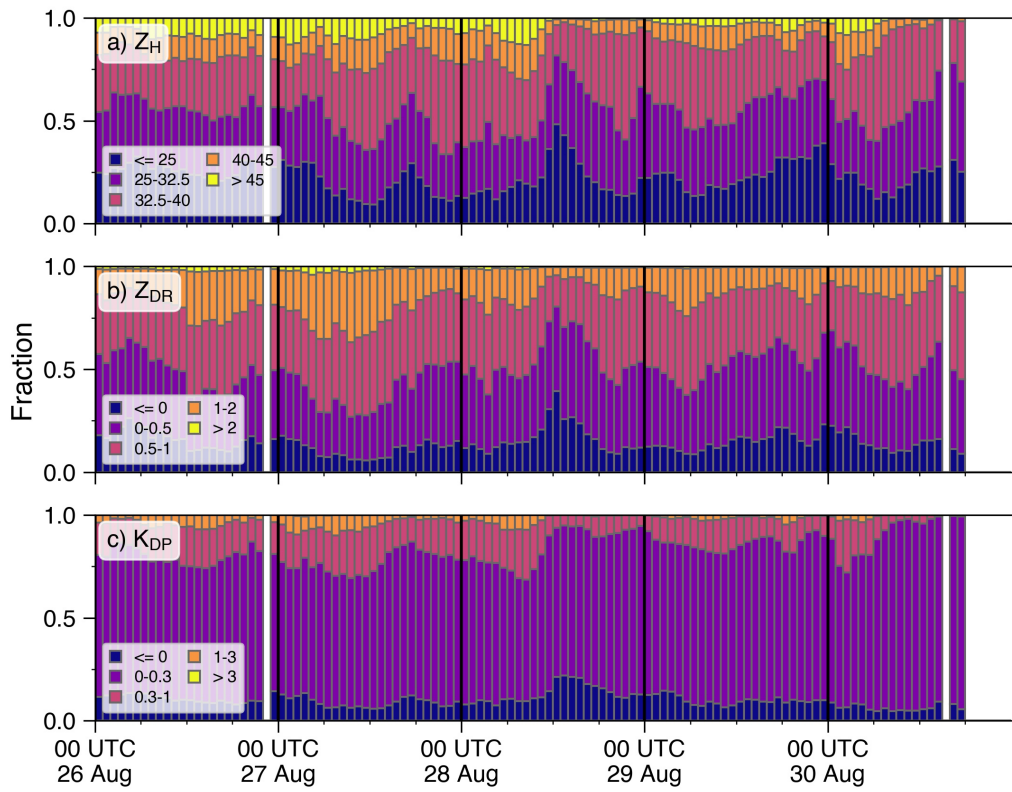




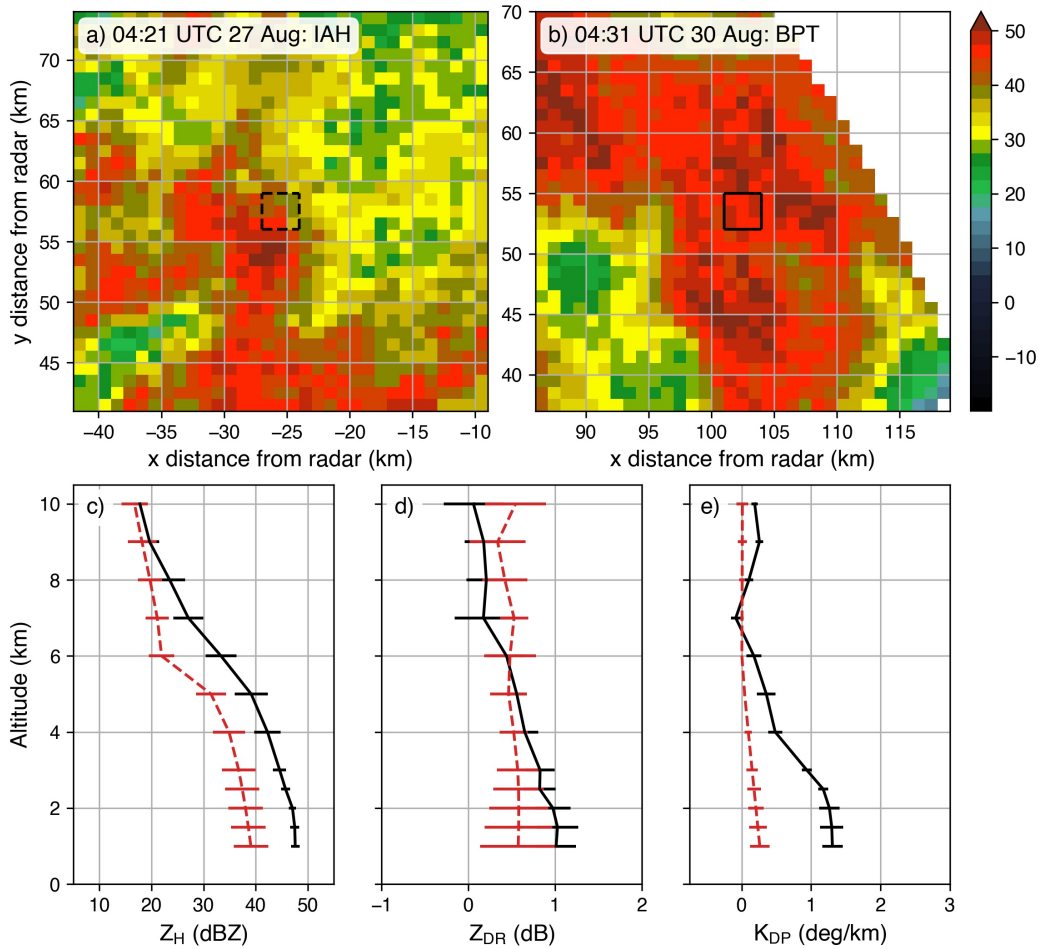
216 **Figure 1.** a) Map of Hurricane Harvey rainfall from 0000 UTC on 25 August to 0000 UTC on 31 August,  
 217 2017. Circle size is proportional to the square root of the total rainfall, to limit the overlap of station plots.  
 218 Within each circle, the lower semicircle (hashed) displays the frequencies and the upper semicircle (solid)  
 219 displays the contribution to the total rainfall by intense (green) and light (white) rainfall. Only hours with  
 220 measurable rainfall are included. Locations marked by an 'x' indicate rain gauges missing more than an av-  
 221 erage of four observations per day. b) Hourly rainfall time series at KIAH. c) Hourly rainfall time series at  
 222 KBPT. Black horizontal line indicates an hourly rain rate of  $25 \text{ mm h}^{-1}$ , which is used to identify intense rain.  
 223 Vertical gray bars indicate missing data.



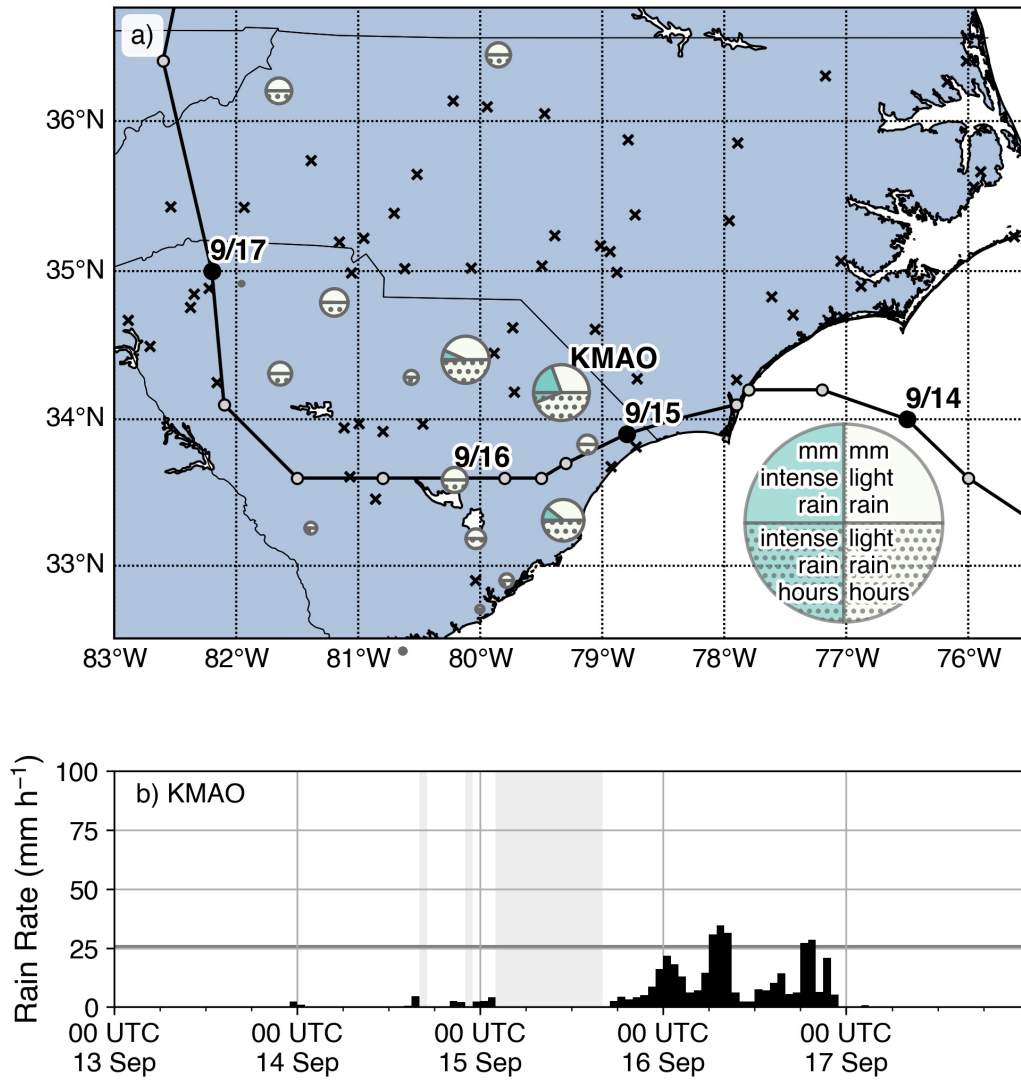
279 **Figure 2.** a) Map of reflectivity from the 0.5° plan position indicator (PPI) scan from KHXG at 0919 UTC  
 280 on 27 August, 2017. Only data within 127 km of KHXG are displayed. The black circles represent the loca-  
 281 tion of ASOS rain gauges within 127 km of KHXG. b) As in a), but for the PID categories. c) Hourly time  
 282 series of range-weighted counts of rain PID categories. d) As in c), but for the range-weighted fraction.



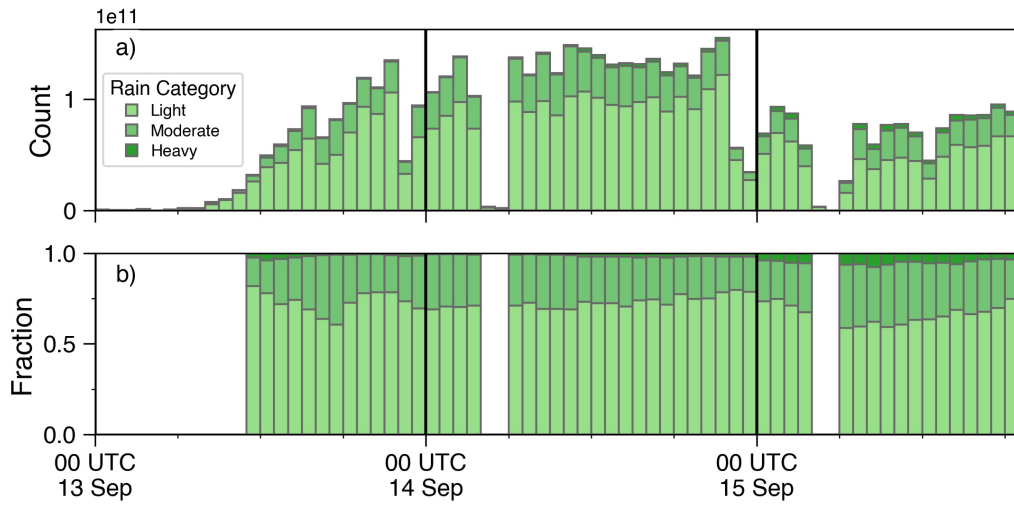
308 **Figure 3.** Hourly time series of range-weighted fraction of polarimetric variables in specified bins at  
 309 KHGX of a) reflectivity (dBZ), b) differential reflectivity (dB), and c) specific differential phase ( $^{\circ} \text{km}^{-1}$ ).



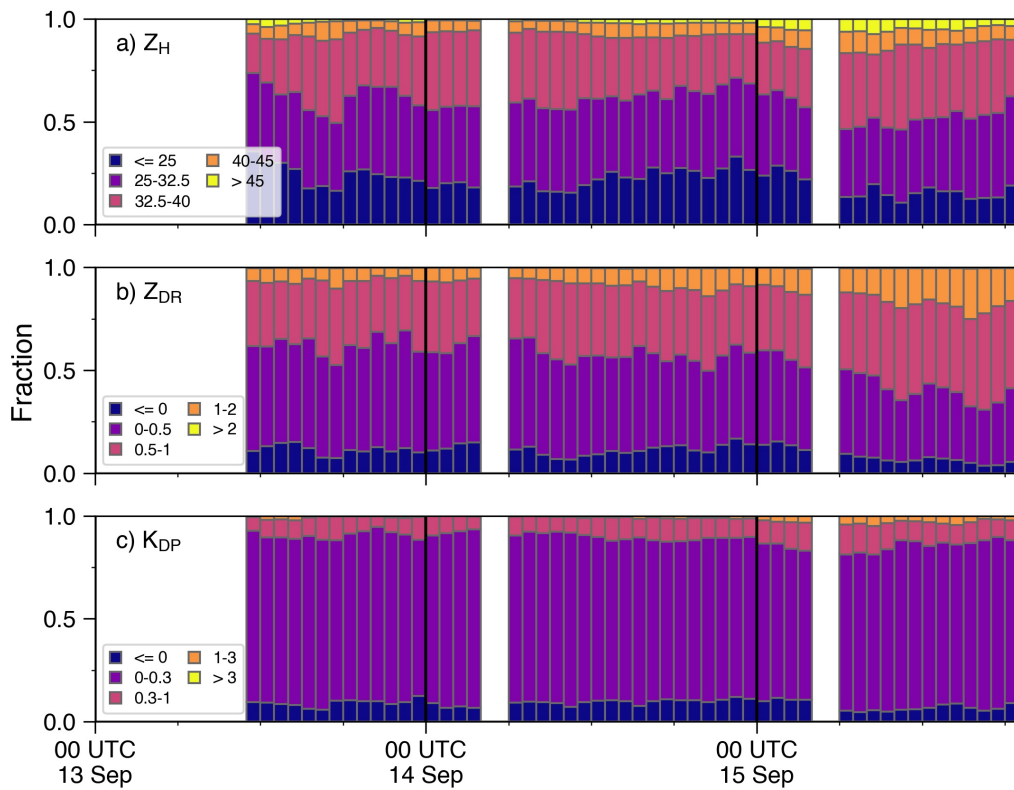
343 **Figure 4.** a) Map of gridded radar reflectivity data at 1 km altitude within a 32x32 km box surrounding  
 344 KIAH at 0421 UTC on 27 August 2017 during Hurricane Harvey. Black dashed box indicates the area over  
 345 which the vertical profiles are calculated. b) As in a), but for KBPT at 0431 UTC on 30 August 2017. Black  
 346 solid box indicates the area over which the vertical profiles are calculated. c) Vertical profiles of  $Z_H$  near  
 347 KIAH (dashed red line) and KBPT (solid black line). Error bars denote the standard deviation at each altitude.  
 348 d) As in c), but for  $Z_{DR}$ . e) As in c), but for  $K_{DP}$



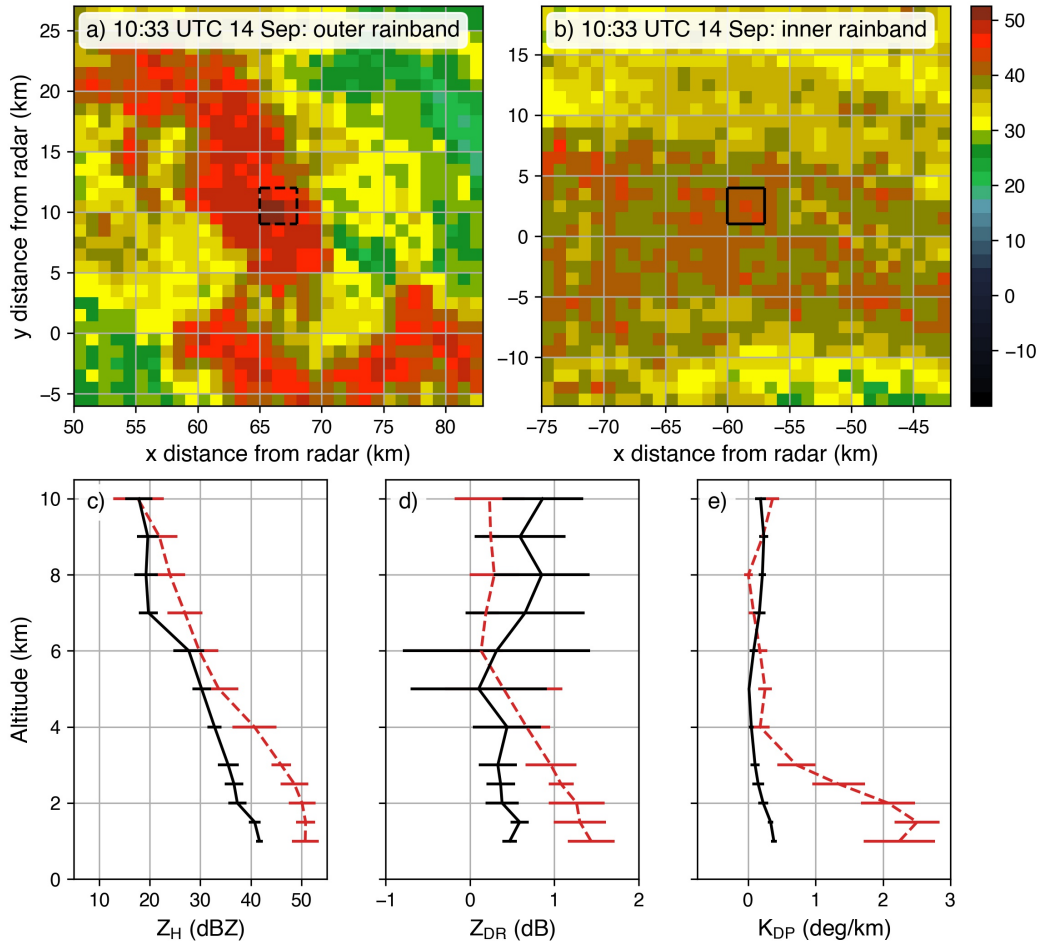
374 **Figure 5.** a) As in Fig. 1, but of Hurricane Florence rainfall from 0000 UTC on 13 September to 0000  
 375 UTC on 18 September, 2018. b) Hourly rainfall time series at KMAO. Black horizontal line indicates an  
 376 hourly rain rate of 25 mm h<sup>-1</sup>, which is used to identify intense rain. Vertical gray bars indicate missing data.



394 **Figure 6.** Hourly time series of range-weighted PID categories from KMHX by a) total count, b) fraction.

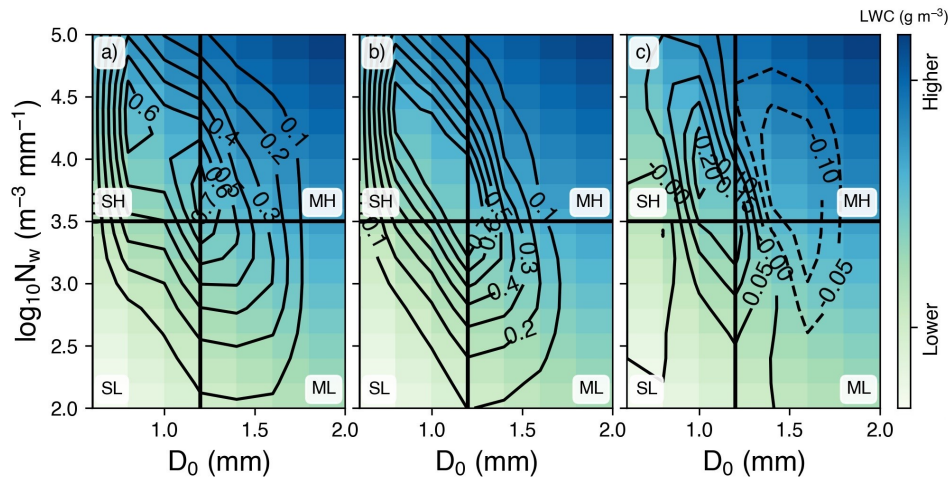


403 **Figure 7.** As in Fig. 3, but for data from KMHX.



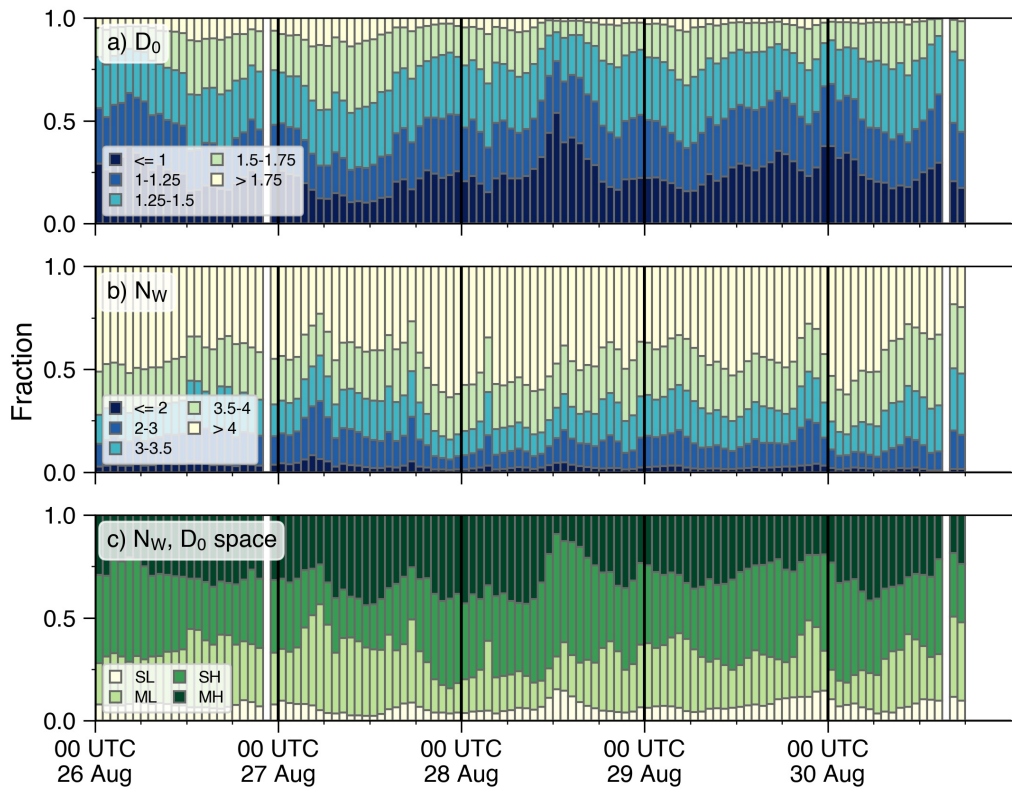
418 **Figure 8.** a) Map of gridded radar reflectivity data at 1 km altitude within a 32x32 km box surrounding an  
 419 outer rainband at 1033 UTC on 14 September 2018 during Hurricane Florence. Black dashed box indicates  
 420 the area over which the vertical profiles are calculated. b) As in a), but surrounding an inner rainband. Black  
 421 solid box indicates the area over which the vertical profiles are calculated. c) Vertical profiles of  $Z_H$  near the  
 422 outer rainband (dashed red line) and inner rainband (solid black line). Error bars denote the standard deviation  
 423 at each altitude. d) As in c), but for  $Z_{DR}$ . e) As in c), but for  $K_{DP}$



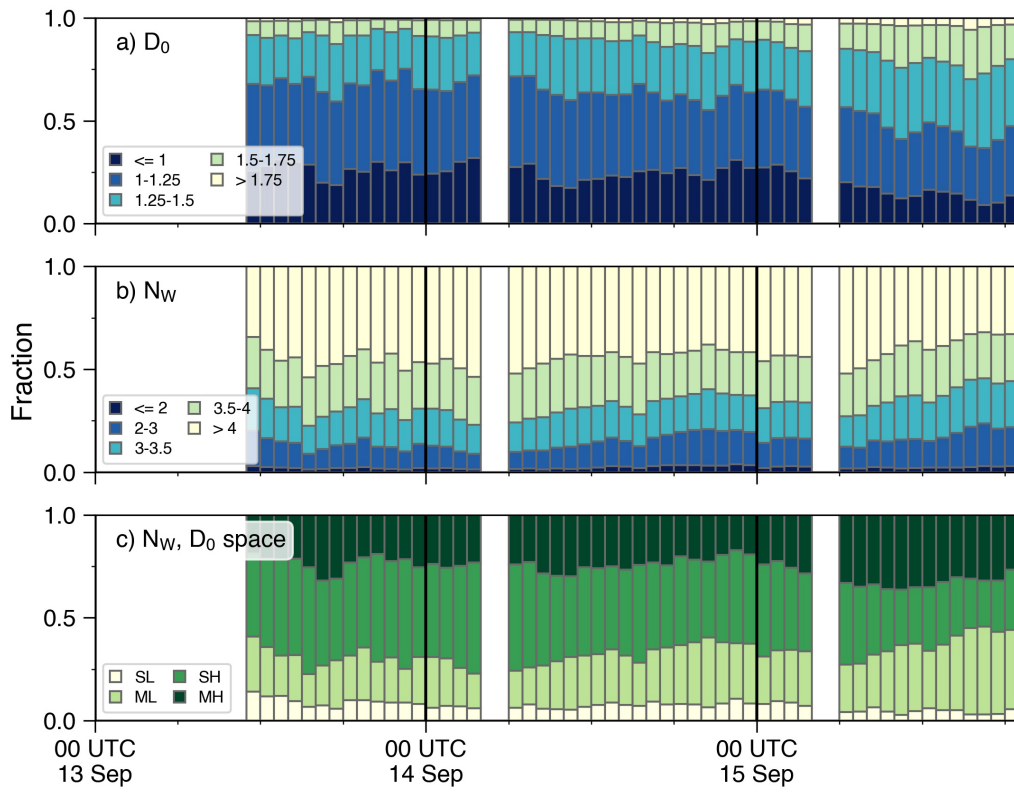


478 **Figure 9.** a) Joint probability distribution (contours) of range-weighted  $N_w$  and  $D_0$  values calculated  
 479 from the KHGX polarimetric data for Harvey between 0000 UTC on 26 August and 0000 UTC on 31 Au-  
 480 gust. Theoretical LWC values (colors) are calculated from Equation 2. Labels indicate the different DSD  
 481 quadrants (SL: low-concentration small drops, ML: low-concentration medium drops, SH: high-concentration  
 482 small drops, and MH: high-concentration medium drops). b) As in a), but for data from KMHX for Florence  
 483 between 0000 UTC on 13 September and 1922 UTC on 15 September, when the radar went offline and c)  
 484 the difference joint probability distribution between Florence/KMHX and Harvey/KHGX (dashed contours  
 485 indicate frequencies were greater for KHGX).

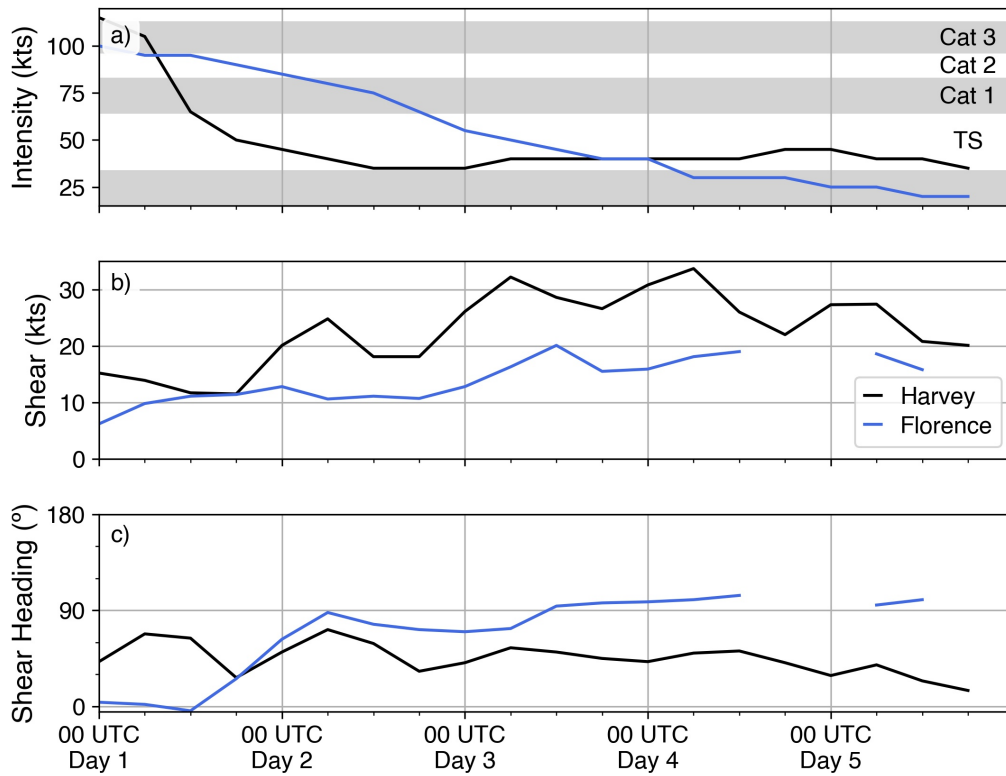




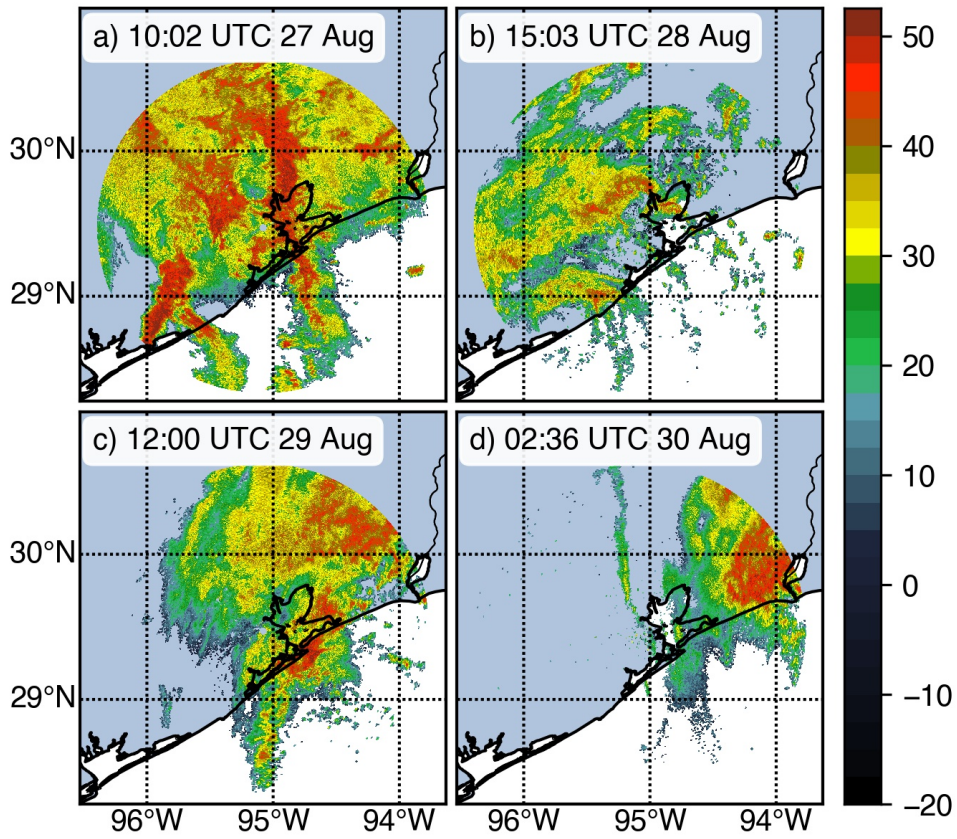
523 **Figure 10.** As in Fig. 3, but for the estimated a)  $D_0$  (mm), b)  $N_W$  ( $\text{m}^{-3} \text{mm}^{-1}$ ), and c)  $N_W, D_0$  quadrants  
 524 defined in Fig. 9 for data from Harvey (2017).



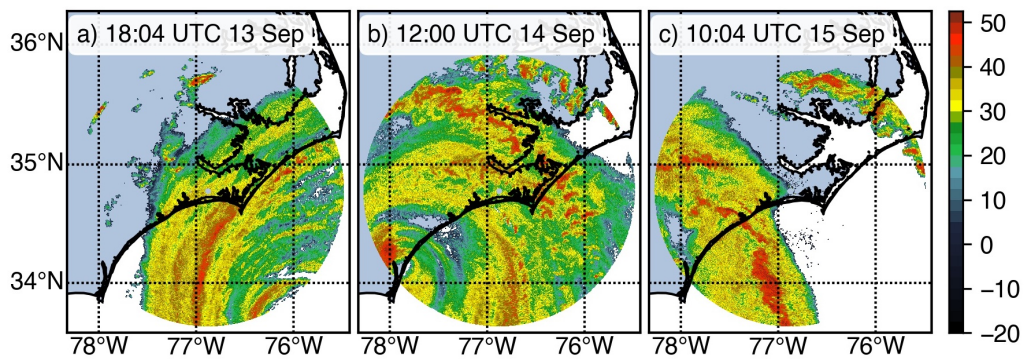
579 **Figure 11.** As in Fig. 7, but for the estimated a)  $D_0$  (mm), b)  $N_W$  ( $\text{m}^{-3} \text{mm}^{-1}$ ), and c)  $N_W, D_0$  quadrants  
 580 defined in Fig. 9 for data from Florence (2018).



582 **Figure 12.** Time series of a) storm intensity (kts), b) deep-layer wind shear magnitude (kts), and c) deep  
 583 layer wind shear heading (°) for Harvey (black) and Florence (blue). Label of Day 1 corresponds to 26 August  
 584 and 13 September for Harvey and Florence, respectively.



596 **Figure 13.** Map of reflectivity from the 0.5°PPI scan from KHXG at a) 1002 UTC on 27 August, b) 1503  
 597 UTC on 28 August, c) 1200 UTC on 29 August, and d) 0236 UTC on 30 August. Only data within 127 km of  
 598 KHXG are displayed.



619 **Figure 14.** Map of reflectivity from the 0.5°PPI scan from KMHX at a) 1804 UTC on 13 September, b)  
620 1200 UTC on 14 September, and c) 1004 UTC on 15 September. Only data within 127 km of KMHX are  
621 displayed. displayed.